# A classification theorem for hypersurfaces of Minkowski spaces 

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#### Abstract

Let $M^{n}$ be a compact hypersurface of a Minkowski space $\left(V^{n+1}, \bar{F}\right)$. In this paper, using the Gauss formula of the Chern connection for Finsler submanifolds, we prove that if the second mean curvature $H_{2}$ of $M$ is constant and the norm square $S$ of the second fundamental form of $M$ satisfies $S \leq \frac{n(n-1)}{n-2} H_{2}$, then $M$ with the induced metric is isometric to the standard Euclidean sphere. This generalizes the result of [2] from the Euclidean to the Minkowski space.


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## 1. Introduction

Let $M$ be an $n$-dimensional smooth manifold and $\pi: T M \rightarrow M$ be the natural projection from the tangent bundle. Let $(x, Y)$ be a point of $T M$ with $x \in M, Y \in$ $T_{x} M$ and let $\left(x^{i}, Y^{i}\right)$ be local coordinates on $T M$ with $Y=Y^{i} \frac{\partial}{\partial x^{i}}$. A Finsler metric on $M$ is a function $F: T M \rightarrow[0,+\infty)$ satisfying the following properties:
(i) Regularity: $F(x, Y)$ is smooth on $T M \backslash 0$;
(ii) Positive homogeneity: $F(x, \lambda Y)=\lambda F(x, Y)$ for $\lambda>0$;
(iii) Strong convexity: The fundamental quadratic form $g_{Y}=g_{i j}(x, Y) d x^{i} \otimes d x^{j}$ is positively definite, where $g_{i j}=\frac{1}{2} \partial^{2}\left(F^{2}\right) / \partial Y^{i} \partial Y^{j}$.

The simplest class of Finsler manifolds is Minkowski space. Let $V^{n+1}$ be a real vector space. A Finsler metric $\bar{F}: T V^{n+1} \rightarrow[0, \infty)$ is called Minkowski if $\bar{F}$ is a function of $\bar{Y} \in V^{n+1}$ only. In this case $\left(V^{n+1}, \bar{F}\right)$ is called a Minkowski space.

Riemannian submanifolds are important in modern differential geometry. There has been a long history for the study of Riemannian submanifolds. For a compact Riemannian hypersurface $M$ of Euclidean space, the second fundamental form is $B=h_{i j}^{n+1} \omega^{i} \otimes \omega^{j} \otimes e_{n+1}$, where $\left\{\omega^{i}\right\}$ is the orthonormal coframe of $M$. The

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second mean curvature $H_{2}$ of $M$ is defined by $H_{2}=\frac{2}{n(n-1)} \sum_{1 \leq i<j \leq n} \lambda_{i} \lambda_{j}$, where $\lambda_{i}$ are the eigenvalues of the second fundamental tensor $h_{i j}^{n+1}$ of $M$. For a compact hypersurface $M$ of Euclidean space, the Gauss equation is $\sum_{i, j} R_{i j i j}=n(n-1) H_{2}$, which implies that the scalar curvature is constant if and only if the second mean curvature $\mathrm{H}_{2}$ is constant.

As well kown, using Cheng-Yau's self-adjoint operator $\square, \mathrm{Li}$ [2] proved that if the second mean curvature $H_{2}$ is constant and the norm square $S$ of the second fundamental form of $M$ satisfies $S \leq \frac{n(n-1)}{n-2} H_{2}$, then $M$ is a Riemannian sphere. As far as we know, there are very few rigidity results on Finsler submanifolds. The main purpose of this paper is to generalize the above result of Li from the Euclidean to the Minkowski space. In this paper, using the Gauss formula for the Chern connection and defining a similar self-adjoint operator $\square$ on Finsler manifolds, we study the hypersurfaces of Minkowski space $\left(V^{n+1}, \bar{F}\right)$ and we obtain the following:
Main Theorem. Let $M^{n}$ be a compact hypersurface of Minkowski space $\left(V^{n+1}, \bar{F}\right)$. If the second mean curvature $H_{2}$ is constant and the norm square $S$ of the second fundamental form of $M$ satisfies $S \leq \frac{n(n-1)}{n-2} H_{2}$, then $M$ with the induced metric is isometric to the standard Euclidean sphere.

## 2. Preliminaries

Let $\left(M^{n}, F\right)$ be an $n$-dimensional Finsler manifold. Then $F$ inherits the Hilbert form, the fundamental tensor and the Cartan tensor as follows [1]:

$$
\begin{aligned}
\omega & =\frac{\partial F}{\partial Y^{i}} d x^{i}, \\
g_{Y} & =g_{i j}(x, Y) d x^{i} \otimes d x^{j}, \\
A_{Y} & =A_{i j k} d x^{i} \otimes d x^{j} \otimes d x^{k}, \\
A_{i j k} & :=\frac{F \partial g_{i j}}{2 \partial Y^{k}} .
\end{aligned}
$$

Let $\varphi:\left(M^{n}, F\right) \rightarrow\left(\bar{M}^{n+p}, \bar{F}\right)$ be an isometric immersion from a Finsler manifold to another one. We have [7]

$$
\begin{gather*}
F(Y)=\bar{F}\left(\varphi_{*}(Y)\right), \quad g_{Y}(U, V)=\bar{g}_{\varphi_{*}(Y)}\left(\varphi_{*}(U), \varphi_{*}(V)\right),  \tag{2.1}\\
A_{Y}(U, V, W)=\bar{A}_{\varphi_{*}(Y)}\left(\varphi_{*}(U), \varphi_{*}(V), \varphi_{*}(W)\right), \tag{2.2}
\end{gather*}
$$

where $Y, \underline{U}, V, W \in T M, \bar{g}$ and $\bar{A}$ are the fundamental tensor and the Cartan tensor of $\bar{M}$, respectively. It can be seen from (2.1) that $\varphi^{*}(\bar{\omega})=\omega$, where $\bar{\omega}$ is the Hilbert form of $\bar{M}$.

In the following we simplify $A_{Y}$ and $g_{Y}$ to $A$ and $g$, respectively. Moreover any vector $U \in T M$ will be identified with the corresponding vector $\varphi_{*}(U) \in T \bar{M}$
and we will use the following convention:

$$
1 \leq i, j, \cdots \leq n, n+1 \leq \alpha, \beta, \cdots \leq n+p, 1 \leq \lambda, \mu, \cdots \leq n-1,1 \leq a, b, \cdots \leq n+p .
$$

Let $\varphi:\left(M^{n}, F\right) \rightarrow\left(\bar{M}^{n+p}, \bar{F}\right)$ be an isometric immersion from a Finsler manifold to another one. Take a $\bar{g}$-orthonormal frame form $\left\{e_{a}\right\}$ for each fibre of $\pi^{*} T \bar{M}$ and let $\left\{\omega^{a}\right\}$ be its local dual coframe, such that $\left\{e_{i}\right\}$ is a frame field for each fibre of $\pi^{*} T M$ and $\omega^{n}$ is the Hilbert form, where $\pi: T M \rightarrow M$ denotes the natural projection. Let $\theta_{b}^{a}$ and $\omega_{j}^{i}$ denote the Chern connection 1-form of $\bar{F}$ and $F$, respectively, i.e. $\bar{\nabla} e_{a}=\theta_{a}^{b} e_{b}$ and $\nabla e_{i}=\omega_{i}^{j} e_{j}$, where $\bar{\nabla}$ and $\nabla$ are the Chern connections of $\bar{M}$ and $M$, respectively. We obtain that $A\left(e_{i}, e_{j}, e_{n}\right)=\bar{A}\left(e_{a}, e_{b}, e_{n}\right)=0$, where $e_{n}=\frac{Y^{i}}{F} \frac{\partial}{\partial x^{i}}$ is the natural dual of the Hilbert form $\omega^{n}$.

The structure equations of $\bar{M}$ are given by

$$
\left\{\begin{array}{l}
d \theta^{a}=-\theta_{b}^{a} \wedge \theta^{b} \\
d \theta_{b}^{a}=-\theta_{c}^{a} \wedge \theta_{b}^{c}+\frac{1}{2} \bar{R}_{b c d}^{a} \omega^{c} \wedge \omega^{d}+\bar{P}_{b c d}^{a} \omega^{c} \wedge \theta_{n}^{d} \\
\theta_{b}^{a}+\theta_{a}^{b}=-2 \bar{A}_{a b c} \theta_{n}^{c}, \\
\theta_{n}^{a}+\theta_{a}^{n}=0, \quad \theta_{n}^{n}=0
\end{array}\right.
$$

By $\theta^{\alpha}=0$ and the structure equations of $\bar{M}$, we have that $\theta_{j}^{\alpha} \wedge \omega^{j}=0$, which implies that $\theta_{j}^{\alpha}=h_{i j}^{\alpha} \omega^{i}, \quad h_{i j}^{\alpha}=h_{j i}^{\alpha}$. We obtain [3]

$$
\begin{equation*}
\omega_{i}^{j}=\theta_{i}^{j}-\Psi_{j i k} \omega^{k} \tag{2.3}
\end{equation*}
$$

where

$$
\begin{align*}
\Psi_{j i k}= & h_{j n}^{\alpha} \bar{A}_{k i \alpha}-h_{k n}^{\alpha} \bar{A}_{j i \alpha}-h_{i n}^{\alpha} \bar{A}_{k j \alpha}-h_{n n}^{\alpha} \bar{A}_{i k s} \bar{A}_{s j \alpha}+h_{n n}^{\alpha} \bar{A}_{i j s} \bar{A}_{s k \alpha}  \tag{2.4}\\
& +h_{n n}^{\alpha} \bar{A}_{j k s} \bar{A}_{s i \alpha} .
\end{align*}
$$

In particular,

$$
\begin{equation*}
\omega_{i}^{n}=\theta_{i}^{n}-h_{n n}^{\alpha} \bar{A}_{k i \alpha} \omega^{k} . \tag{2.5}
\end{equation*}
$$

Using the almost $\bar{g}$-compatibility, we have

$$
\begin{equation*}
\theta_{\alpha}^{j}=\left(-h_{i j}^{\alpha}-2 h_{n i}^{\beta} \bar{A}_{j \alpha \beta}+2 h_{n n}^{\beta} \bar{A}_{j \lambda \alpha} \bar{A}_{i \lambda \beta}\right) \omega^{i}-2 \bar{A}_{j \alpha \lambda} \omega_{n}^{\lambda} . \tag{2.6}
\end{equation*}
$$

In particular, $\theta_{\alpha}^{n}=-h_{n i}^{\alpha} \omega^{i}$. We quote the following results:
Proposition 2.1 ([3]). Let $\varphi:\left(M^{n}, F\right) \rightarrow\left(\bar{M}^{n+p}, \bar{F}\right)$ be an isometric immersion from a Finsler manifold to a Minkowski space. Then

$$
\left\{\begin{array}{l}
\left(\bar{\nabla}_{e_{i}} \bar{A}\right)(\bullet, \bullet, \bullet)=0, \\
\bar{A}\left(\bullet, \bullet, \bar{\nabla}_{e_{i}} e_{n}\right)=0
\end{array}\right.
$$

Proposition 2.2 ([3] Gauss equations). Let $\varphi:\left(M^{n}, F\right) \rightarrow\left(\bar{M}^{n+p}, \bar{F}\right)$ be an isometric immersion from a Finsler manifold another one. Then we have

$$
\left\{\begin{aligned}
P_{i k \lambda}^{j}= & \bar{P}_{i k \lambda}^{j}+\Psi_{j i k ; \lambda}-2 \Psi_{s i k} A_{j s \lambda}-2 h_{i k}^{\alpha} \bar{A}_{j \lambda \alpha} \\
R_{i k l}^{j}= & \bar{R}_{i k l}^{j}-h_{i k}^{\alpha} h_{j l}^{\alpha}+h_{i l}^{\alpha} h_{j k}^{\alpha}+\Psi_{j i k \mid l}-\Psi_{j i l \mid k} \\
& +\Psi_{s i k} \Psi_{j s l}-\Psi_{s i l} \Psi_{j s k}-2 h_{i k}^{\alpha} h_{n l}^{\beta} \bar{A}_{j \alpha \beta}+2 h_{i l}^{\alpha} h_{n k}^{\beta} \bar{A}_{j \alpha \beta} \\
& +2 h_{i k}^{\alpha} h_{n n}^{\beta} \bar{A}_{j s \alpha} \bar{A}_{l s \beta}-2 h_{i l}^{\alpha} h_{n n}^{\beta} \bar{A}_{j s \alpha} \bar{A}_{k s \beta}-h_{n n}^{\alpha} \bar{A}_{s l \alpha} \bar{P}_{i k s}^{j} \\
& +h_{n n}^{\alpha} \bar{A}_{s k \alpha} \bar{P}_{i l s}^{j}+h_{n l}^{\alpha} \bar{P}_{i k \alpha}^{j}-h_{n k}^{\alpha} \bar{P}_{i l \alpha}^{j}
\end{aligned}\right.
$$

where ";" and "|" respectively denote the vertical and the horizontal covariant differentials with respect to the Chern connection $\nabla$.
Proposition 2.3 ([3] Codazzi equations). Let $\varphi:\left(M^{n}, F\right) \rightarrow\left(\bar{M}^{n+p}, \bar{F}\right)$ be an isometric immersion from a Finsler manifold to a Finsler manifold to another one. Then we have

$$
\left\{\begin{array}{l}
h_{i j ; \lambda}^{\alpha}=-\bar{P}_{i j \lambda}^{\alpha} \\
h_{i j \mid k}^{\alpha}-h_{i k \mid j}^{\alpha}=-\bar{R}_{i j k}^{\alpha}+h_{n j}^{\beta} \bar{P}_{i k \beta}^{\alpha}-h_{n k}^{\beta} \bar{P}_{i j \beta}^{\alpha} \\
-h_{l k}^{\alpha} \Psi_{l i j}+h_{l j}^{\alpha} \Psi_{l i k}-h_{n n}^{\beta} \bar{A}_{l j \beta} \bar{P}_{i k l}^{\alpha}+h_{n n}^{\beta} \bar{A}_{l k \beta} \bar{P}_{i j l}^{\alpha}
\end{array}\right.
$$

## 3. Hypersurfaces of a Minkowski space

Let $\left(M^{n}, F\right)$ be a compact hypersurfaces of a Minkowski space $\left(V^{n+1}, \bar{F}\right)$. Then we have

$$
\begin{equation*}
h_{i j \mid k}^{n+1} \omega^{k}+h_{i j ; \lambda}^{n+1} \omega_{n}^{\lambda}=d h_{i j}^{n+1}-h_{k j}^{n+1} \omega_{i}^{k}-h_{i k}^{n+1} \omega_{j}^{k}+h_{i j}^{n+1} \theta_{n+1}^{n+1} \tag{3.1}
\end{equation*}
$$

Exterior differentiation of the left-hand side terms of (3.1), gives

$$
\begin{align*}
d & h_{i j \mid k}^{n+1} \wedge \omega^{k}+h_{i j \mid k}^{n+1} d \omega^{k}+d h_{i j ; \lambda}^{n+1} \wedge \omega_{n}^{\lambda}+h_{i j ; \lambda}^{n+1} d \omega_{n}^{\lambda} \\
= & \left\{h_{i j|k| l}^{n+1} \omega^{l}+h_{i j \mid k ; \mu}^{n+1} \omega_{n}^{\mu}+h_{l j \mid k}^{n+1} \omega_{i}^{l}+h_{i l \mid k}^{n+1} \omega_{j}^{l}+h_{i j \mid l}^{n+1} \omega_{k}^{l}-h_{i j \mid k}^{n+1} \theta_{n+1}^{n+1}\right\} \wedge \omega^{k} \\
& +h_{i j \mid k}^{n+1}\left\{-\omega_{l}^{k} \wedge \omega^{l}\right\} \\
& +\left\{h_{i j ; \lambda \mid l}^{n+1} \omega^{l}+h_{i j ; \lambda ; \mu}^{n+1} \omega_{n}^{\mu}+h_{l j ; \lambda}^{n+1} \omega_{i}^{l}+h_{i l ; \lambda}^{n+1} \omega_{j}^{l}+h_{i j ; \mu}^{n+1} \omega_{\lambda}^{\mu}-h_{i j ; \lambda}^{n+1} \theta_{n+1}^{n+1}\right\} \wedge \omega_{n}^{\lambda} \\
& +h_{i j ; \lambda}^{n+1}\left\{-\omega_{\mu}^{\lambda} \wedge \omega_{n}^{\mu}+\frac{1}{2} R_{n l s}^{\lambda} \omega^{l} \wedge \omega^{s}+P_{n l \mu}^{\lambda} \omega^{l} \wedge \omega_{n}^{\mu}\right\}  \tag{3.2}\\
= & \left\{-h_{i j|k| l}^{n+1}+\frac{1}{2} h_{i j ; \lambda}^{n+1} R_{n k l}^{\lambda}\right\} \omega^{k} \wedge \omega^{l}-h_{i j ; \lambda ; \mu}^{n+1} \omega_{n}^{\lambda} \wedge \omega_{n}^{\mu} \\
& +\left\{-h_{i j \mid k ; \lambda}^{n+1}+h_{i j ; \lambda \mid k}^{n+1}+h_{i j ; \mu}^{n+1} P_{n k l}^{\mu}+h_{i j \mid k}^{n+1} \bar{A}_{n+1 n+1 \lambda}\right\} \omega^{k} \wedge \omega_{n}^{\lambda} \\
& +h_{l j \mid k}^{n+1} \omega_{i}^{l} \wedge \omega^{k}+h_{i l \mid k}^{n+1} \omega_{j}^{l} \wedge \omega^{k}+h_{l j ; \lambda}^{n+1} \omega_{i}^{l} \wedge \omega_{n}^{\lambda}+h_{i l ; \lambda}^{n+1} \omega_{j}^{l} \wedge \omega_{n}^{\lambda} \\
& -h_{i j \mid k}^{n+1} \theta_{n+1}^{n+1} \wedge \omega^{k}-h_{i j ; \lambda}^{n+1} \theta_{n+1}^{n+1} \wedge \omega_{n}^{\lambda} .
\end{align*}
$$

Exterior differentiation of the right-hand side terms of (3.1), gives also

$$
\begin{align*}
& -d h_{k j}^{n+1} \wedge \omega_{i}^{k}-h_{k j}^{n+1} d \omega_{i}^{k}-d h_{i k}^{n+1} \wedge \omega_{j}^{k}-h_{i k}^{n+1} d \omega_{j}^{k} \\
& +d h_{i j}^{n+1} \wedge \theta_{n+1}^{n+1}+h_{i j}^{n+1} d \theta_{n+1}^{n+1} \\
= & -\left\{h_{k j \mid l}^{n+1} \omega^{l}+h_{k j ; \lambda}^{n+1} \omega_{n}^{\lambda}+h_{l j}^{n+1} \omega_{k}^{l}+h_{k l}^{n+1} \omega_{j}^{l}-h_{k j}^{n+1} \theta_{n+1}^{n+1}\right\} \wedge \omega_{i}^{k} \\
& -h_{k j}^{n+1}\left\{-\omega_{l}^{k} \wedge \omega_{i}^{l}+\frac{1}{2} R_{i l s}^{k} \omega^{l} \wedge \omega^{s}+P_{i l \lambda}^{k} \omega^{l} \wedge \omega_{n}^{\lambda}\right\} \\
& -\left\{h_{i k \mid l}^{n+1} \omega^{l}+h_{i k ; \lambda}^{n+1} \omega_{n}^{\lambda}+h_{l k}^{n+1} \omega_{i}^{l}+h_{i l}^{n+1} \omega_{k}^{l}-h_{i k}^{n+1} \theta_{n+1}^{n+1}\right\} \wedge \omega_{j}^{k} \\
& -h_{i k}^{n+1}\left\{-\omega_{l}^{k} \wedge \omega_{j}^{l}+\frac{1}{2} R_{j l s}^{k} \omega^{l} \wedge \omega^{s}+P_{j l \lambda}^{k} \omega^{l} \wedge \omega_{n}^{\lambda}\right\} \\
& +\left\{h_{i j \mid k}^{n+1} \omega^{k}+h_{i j ; \lambda}^{n+1} \omega_{n}^{\lambda}+h_{k j}^{n+1} \omega_{i}^{k}+h_{i k}^{n+1} \omega_{j}^{k}\right\} \wedge \theta_{n+1}^{n+1} \\
& +h_{i j}^{n+1}\left\{2 h_{s k}^{n+1} h_{n l}^{n+1} \bar{A}_{s n+1 n+1}-2 h_{s k}^{n+1} h_{n n}^{n+1} \bar{A}_{s t n+1} \bar{A}_{t l n+1}\right\} \omega^{k} \wedge \omega^{l}  \tag{3.3}\\
& +2 h_{i j}^{n+1} h_{s k}^{n+1} \bar{A}_{s n n+\lambda} \omega^{k} \wedge \omega_{n}^{\lambda} \\
= & -\frac{1}{2} h_{s j}^{n+1} R_{i k l}^{s}-\frac{1}{2} h_{i s}^{n+1} R_{j k l}^{k}+2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n l}^{n+1} \bar{A}_{s n+1 n+1} \\
& \left.-2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n n}^{n+1} \bar{A}_{s t n+1} \bar{A}_{t l n}+1\right\} \omega^{k} \wedge \omega^{l} \\
& +\left\{-h_{s j}^{n+1} P_{i k \lambda}^{s}-h_{s j}^{n+1} P_{j k \lambda}^{k}+2 h_{i j}^{n+1} h_{s k}^{n+1} \bar{A}_{s n+1 \lambda}\right\} \omega^{k} \wedge \omega_{n}^{\lambda} \\
& -h_{k j \mid l}^{n+1} \omega^{l} \wedge \omega_{i}^{k}-h_{k j ; \lambda}^{n+1} \omega_{n}^{\lambda} \wedge \omega_{i}^{k}-h_{i k \mid l}^{n+1} \omega^{l} \wedge \omega_{j}^{k}-h_{i k ; \lambda}^{n+1} \omega_{n}^{\lambda} \wedge \omega_{j}^{k} \\
& -h_{i j \mid k}^{n+1} \theta_{n+1}^{n+1} \wedge \omega^{k}-h_{i j ; \lambda}^{n+1} \theta_{n+1}^{n+1} \wedge \omega_{n}^{\lambda} .
\end{align*}
$$

It can be seen from (3.2) and (3.3) that

$$
\begin{align*}
& \left\{h_{i j|k| l}^{n+1}-\frac{1}{2} h_{s j}^{n+1} R_{i k l}^{s}-\frac{1}{2} h_{i s}^{n+1} R_{j k l}^{s}-\frac{1}{2} h_{i j ; \lambda}^{n+1} R_{n k l}^{\lambda}\right. \\
& \left.\quad+2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n l}^{n+1} \bar{A}_{s n+1 n+1}-2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n n}^{n+1} \bar{A}_{s t n+1} \bar{A}_{t l n+1}\right\} \omega^{k} \wedge \omega^{l} \\
& +  \tag{3.4}\\
& +\left\{h_{i j \mid k ; \lambda}^{n+1}-h_{i j ; \lambda \mid k}^{n+1}+h_{i j ; \mu}^{n+1} P_{n k \lambda}^{\mu}+h_{i s}^{n+1} P_{j k l}^{s}+h_{s j}^{n+1} P_{i k l}^{s}-h_{i j \mid k}^{n+1} \bar{A}_{n+1 n+1 \lambda}\right. \\
& \left.\quad+2 h_{i j}^{n+1} h_{s k}^{n+1} \bar{A}_{s n+1 \lambda}\right\} \omega^{k} \wedge \omega_{n}^{\lambda}+h_{i j ; \lambda ; \mu}^{n+1} \omega_{n}^{\lambda} \wedge \omega_{n}^{\mu} \\
& =
\end{align*}
$$

From (3.4) we immediately obtain the following:
Proposition 3.1. If $M^{n}$ be a hypersurface of Minkowski space $\left(V^{n+1}, \bar{F}\right)$, then

$$
\left\{\begin{aligned}
h_{i j ; \lambda ; \mu}^{n+1}-h_{i j ; \mu ; \lambda}^{n+1}= & 0 \\
h_{i j \mid k ; \lambda}^{n+1}-h_{i j ; \lambda \mid k}^{n+1}= & -h_{s j}^{n+1} P_{i k l}^{s}-h_{i s}^{n+1} P_{j k l}^{s}-h_{i j ; \mu}^{n+1} P_{n k \lambda}^{\mu}+h_{i j \mid k}^{n+1} \bar{A}_{n+1 n+1 \lambda} \\
& -2 h_{i j}^{n+1} h_{s k}^{n+1} \bar{A}_{s n+1 \lambda}, \\
h_{i j|k| l}^{n+1}-h_{i j|l| k}^{n+1}= & h_{s j}^{n+1} R_{i k l}^{s}+h_{i s}^{n+1} R_{j k l}^{s}+h_{i j ; \lambda}^{n+1} R_{n k l}^{\lambda} \\
& -2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n l}^{n+1} \bar{A}_{s n+1 n+1}+2 h_{i j}^{n+1} h_{s l}^{n+1} h_{n k}^{n+1} \bar{A}_{s n+1 n+1} \\
& +2 h_{i j}^{n+1} h_{s k}^{n+1} h_{n n}^{n+1} \bar{A}_{s t n+1} \bar{A}_{t l n+1} \\
& -2 h_{i j}^{n+1} h_{s l}^{n+1} h_{n n}^{n+1} \bar{A}_{s t n+1} \bar{A}_{t k n+1} .
\end{aligned}\right.
$$

The form $B=h_{i j}^{n+1} \omega^{i} \otimes \omega^{j} \otimes e_{n+1}$ is called the second fundamental form of $M$ and $H=\frac{1}{n} \operatorname{tr} B=\frac{1}{n} \sum_{i} h_{i i}^{n+1} e_{n+1}$ is called the mean curvature vector. The norm square $S$ of the second fundamental form of $M$ is $S=\sum_{i j}\left(h_{i j}^{n+1}\right)^{2}$. Let $\lambda_{i}$ be the eigenvalues of the second fundamental tensor $h_{i j}^{n+1}$ of $M$. The second mean curvature $H_{2}=\frac{2}{n(n-1)} \sum_{1 \leq i<j \leq n} \lambda_{i} \lambda_{j}=\frac{1}{n(n-1)}\left[(n H)^{2}-S\right]$.

From the second formula of Proposition 2.1, we obtain that

$$
\begin{equation*}
\bar{A}\left(\bullet, \bullet, \nabla_{e_{i}} e_{n}\right)+\bar{A}\left(\bullet, \bullet, e_{\lambda}\right) \Psi_{\lambda n i}+\bar{A}\left(\bullet, \bullet, e_{n+1}\right) h_{n i}^{n+1}=0 \tag{3.6}
\end{equation*}
$$

Let $P$ be an arbitrary point in $M$. There exists a local coordinate system $\left\{x^{i}\right\}$ such that $\frac{Y^{n}}{F} \frac{\partial}{\partial x^{n}}=e_{n}$. Let $\gamma(t)$ be a curve in $M$ with $\gamma(0)=P$ and tangent vector field $\dot{\gamma}(t)=e_{n}$. Let $X_{i}(t)$ be parallel vector fields along $\gamma(t)$ with $X_{i}(0)=\left.\frac{\partial}{\partial x^{i}}\right|_{p}$. We have that $\nabla_{e_{n}} \frac{\partial}{\partial x^{i}}=\frac{Y^{n}}{F} \nabla_{\frac{\partial}{\partial x^{n}}} \frac{\partial}{\partial x^{i}}=0$ at $P$, i.e. $\Gamma_{n i}^{k}=0$ at $P$, then we can obtain that at $P$

$$
\begin{align*}
\bar{A}\left(\bullet, \bullet, \nabla_{e_{i}} e_{n}\right) & =\bar{A}\left(\bullet, \bullet, \frac{Y^{n}}{F} \nabla_{u_{i}^{j}} \frac{\partial}{\partial x^{j}} \frac{\partial}{\partial x^{n}}\right)  \tag{3.7}\\
& =\bar{A}\left(\bullet, \bullet, \frac{Y^{n}}{F} u_{i}^{j} \Gamma_{j n}^{k} \frac{\partial}{\partial x^{k}}\right)=0 .
\end{align*}
$$

Substituting (3.7) into (3.6) yields that $\bar{A}\left(\bullet, \bullet, e_{j}\right) \Psi_{j n i}+\bar{A}\left(\bullet, \bullet, e_{n+1}\right) h_{n i}^{n+1}=0$ at $P$, which together with (2.4) yields that $-h_{n n}^{n+1} \bar{A}_{j i n+1} \bar{A}_{s t j}+h_{n i}^{n+1} \bar{A}_{s t n+1}=0$ at $P$, so by (2.4) we have that

$$
\begin{equation*}
\Psi_{i j k}=0 \text { and } h_{n i}^{n+1} \bar{A}\left(\bullet, \bullet, e_{n+1}\right)=0, \quad \forall i, j, k, \text { at } P \tag{3.8}
\end{equation*}
$$

It follows from the first formula of Proposition 2.1 and $\bar{A}(\bullet, \bullet, \bullet)_{\mid i}=0$ that

$$
\begin{equation*}
\bar{A}(\bullet, \bullet, \bullet)_{; \lambda} \Psi_{\lambda n i}+\bar{A}(\bullet, \bullet, \bullet)_{; n+1} h_{n i}^{n+1}=0 \text { at } P \tag{3.9}
\end{equation*}
$$

It can be seen from (3.8) and (3.9) that

$$
\begin{equation*}
\bar{A}(\bullet, \bullet, \bullet)_{; n+1} h_{n i}^{n+1}=0 \text { at } P \tag{3.10}
\end{equation*}
$$

Now taking the exterior differentiation of $A_{i j k}=\bar{A}_{i j k}$, we obtain that

$$
\begin{align*}
A_{i j k \mid l}= & \bar{A}_{i j k ; \lambda} \Psi_{\lambda n l}+\bar{A}_{i j k ; n+1} h_{n l}^{n+1} \\
& +\bar{A}_{s j k} \Psi_{s i l}+\bar{A}_{i s k} \Psi_{s j l}+\bar{A}_{i j s} \Psi_{s k l}  \tag{3.11}\\
& +\bar{A}_{n+1 j k} h_{i l}^{n+1}+\bar{A}_{i n+1 k} h_{j l}^{n+1}+\bar{A}_{i j n+1} h_{k l}^{n+1}
\end{align*}
$$

By (3.8) and (3.11), we obtain that

$$
\begin{equation*}
A_{i j k \mid n}=0 \quad \text { at } P \tag{3.12}
\end{equation*}
$$

Define $\delta Y^{i}=d Y^{i}+N_{j}^{i} d x^{j}$. The pull-back of the Sasaki metric $g_{i j} d x^{i} \otimes d x^{j}+$ $g_{i j} \delta Y^{i} \otimes \delta Y^{j}$ from $T M \backslash\{0\}$ to the sphere bundle $S M$ is a Riemannian metric $\widehat{g}=$ $g_{i j} d x^{i} \otimes d x^{j}+\delta_{a b} \omega_{n}^{a} \otimes \omega_{n}^{b}$.

We quote the following results:
Lemma 3.2 ([5]). For $X=\sum_{i} x_{i} \omega^{i} \in \Gamma\left(\pi^{*} T^{*} M\right), \operatorname{div}_{\widehat{g}} X=\sum_{i} x_{i \mid i}+\sum_{\mu, \lambda} x_{\mu} P_{\lambda \lambda \mu}^{n}$.
Lemma 3.3 ([6]). Let $B$ be a real symmetric matrix with $\operatorname{tr} B=0$. Then

$$
\left|\operatorname{tr} B^{3}\right| \leq \frac{n-2}{\sqrt{n(n-1)}}\left(\operatorname{tr} B^{2}\right)^{\frac{3}{2}}
$$

Lemma 3.4 ([4]). All Landsberg spaces of nonzero constant flag curvature must be Riemannian.

Let $\phi=\sum_{i, j} \phi_{i j} \omega^{i} \otimes \omega^{j}$ be a symmetric tensor defined on the sphere bundle $S M$ and $\psi=\sum_{i} \psi_{i} \omega^{i} \in \Gamma\left(\pi^{*} T^{*} M\right)$. Now we can define an operator $\square$ associated to $\phi$ by

$$
\begin{equation*}
\square f=\sum_{i, j} \phi_{i j} f_{|i| j}+\sum_{i, \lambda, \mu} \phi_{i \lambda} f_{\mid i} P_{\mu \mu \lambda}^{n}+\sum_{i} \psi_{i} f_{\mid i}, \quad \forall f \in C^{\infty}(S M) \tag{3.13}
\end{equation*}
$$

Proposition 3.5. Let $(M, F)$ be a compact manifold. Then the operatoris selfadjoint if and only if $\sum_{j} \phi_{i j \mid j}-\psi_{i}=0$.
Proof. Let $X=\sum_{i, j} g \phi_{i j} f_{\mid i} \omega^{j} \in \Gamma\left(\pi^{*} T^{*} M\right), \quad \forall f, g \in C^{\infty}(S M)$. Then we have from Lemma 3.2

$$
\begin{align*}
\operatorname{div}_{\widehat{g}} X & =\sum_{i, j}\left\{g_{\mid j} \phi_{i j} f_{\mid i}+g \phi_{i j \mid j} f_{\mid i}+g \phi_{i j} f_{|i| j}\right\}+\sum_{i, \lambda, \mu} g \phi_{i \lambda} f_{\mid i} P_{\mu \mu \lambda}^{n} \\
& =\sum_{i, j} g_{\mid j} \phi_{i j} f_{\mid i}+\sum_{i} g \psi_{i} f_{\mid i}+\sum_{i, j} g \phi_{i j} f_{|i| j}+\sum_{i, \lambda, \mu} g \phi_{i \lambda} f_{\mid i} P_{\mu \mu \lambda}^{n} \tag{3.14}
\end{align*}
$$

Integrating (3.14) yields

$$
\begin{equation*}
\int_{S M}(\square f) g d V_{S M}=-\int_{S M} \sum_{i, j} g_{\mid j} \phi_{i j} f_{\mid i} d V_{S M} \tag{3.15}
\end{equation*}
$$

Hence we get

$$
\begin{equation*}
\int_{S M}(\square f) g d V_{S M}=\int_{S M}(\square g) f d V_{S M}, \tag{3.16}
\end{equation*}
$$

which completes the proof.

## 4. Proof of the main theorem

Substituting (3.8) into the second formula of Proposition 2.2, by Proposition 2.3 and Proposition 3.1, we obtain that at $P$

$$
\left\{\begin{array}{l}
h_{i j|k| l}^{n+1}=h_{i k|j| l}^{n+1}+h_{s k}^{n+1} \Psi_{s i j \mid l}+h_{s j}^{n+1} \Psi_{s i k \mid l}  \tag{4.1}\\
h_{i j j|k| l}^{n+1}=h_{i j j|l| k}^{n+1}+h_{s j}^{n+1} R_{i k l}^{s}+h_{i s}^{n+1} R_{j k l}^{s} \\
R_{i k l}^{j}=-h_{i k}^{\alpha} h_{j l}^{\alpha}+h_{i l}^{\alpha} h_{j k}^{\alpha}+\Psi_{j i k \mid l}-\Psi_{j i l \mid k}
\end{array}\right.
$$

Let $\omega=d S=S_{\mid i} \omega^{i}+S_{; i} \omega_{n}^{i}$. Then $\omega$ is a global section of $\pi^{*} T^{*} M$. By the first formula of (2.24), we have $S_{; i}=0$, i.e. $\omega=d S=S_{\mid i} \omega^{i}$. In the following, the computation is pointwisely estimated. Using the first formula of (4.1) and Lemma 3.2 , we have that

$$
\begin{align*}
\operatorname{div} \hat{g} \omega= & 2\left[\sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid k}^{n+1}\right]_{\mid k}+2 \sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid \lambda}^{n+1} P_{\mu \mu \lambda}^{n} \\
= & 2 \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}+2 \sum_{i, j, k} h_{i j}^{n+1} h_{i k|j| k}^{n+1}+2 \sum_{i, j, k, s} h_{i j}^{n+1} h_{s k}^{n+1} \Psi_{s i j \mid k}  \tag{4.2}\\
& +2 \sum_{i, j, k, s} h_{i j}^{n+1} h_{s j}^{n+1} \Psi_{s i k \mid k}+2 \sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid \lambda}^{n+1} P_{\mu \mu \lambda}^{n}
\end{align*}
$$

It can be seen from (4.1) and (4.2) that

$$
\begin{align*}
\operatorname{div} \omega= & 2 \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}+2 \sum_{i, j, k, s} h_{i j}^{n+1}\left\{h_{k i|k| j}^{n+1}+h_{s i}^{n+1} R_{k j k}^{s}+h_{k s}^{n+1} R_{i j k}^{s}\right\} \\
& +2 \sum_{i, j, k, s} h_{i j}^{n+} h_{s k}^{n+1} \Psi_{s i j \mid k}+2 \sum_{i, j, k, s} h_{i j}^{n+1} h_{s j}^{n+1} \Psi_{s i k \mid k}+2 \sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid \lambda}^{n+1} P_{\mu \mu \lambda}^{n} \\
= & \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}+2 \sum_{i, j, k, s} h_{i j}^{n+1} h_{k k|i| j}^{n+1}+2 \sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid \lambda}^{n+1} P_{\mu \mu \lambda}^{n} \\
& +n H \sum_{i}\left(\lambda_{i}-H\right)^{3}+3 n H^{2} S-2 n^{2} H^{4}-S^{2}  \tag{4.3}\\
\geq & \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}+\sum_{i, j, k, s} h_{i j}^{n+1} h_{k k|i| j}^{n+1}+\sum_{i, j, k} h_{i j}^{n+1} h_{i j \mid \lambda}^{n+1} P_{\mu \mu \lambda}^{n} \\
& +\frac{n-1}{n}\left[S-n H_{2}\right]\left\{2(n-1) H_{2}-\frac{n-2}{n} S\right. \\
& \left.-\frac{n-2}{n} \sqrt{\left(n(n-1) H_{2}+S\right)\left(S-n H_{2}\right)}\right\}
\end{align*}
$$

where the second mean curvature $H_{2}=\frac{1}{n(n-1)}\left[(n H)^{2}-S\right]$. Let

$$
\begin{aligned}
\square f= & \sum_{i, j}\left(n H \delta_{i j}-h_{i j}^{n+1}\right) f_{|i| j}+\sum_{i, \lambda, \mu}\left(n H \delta_{i \lambda}-h_{i \lambda}^{n+1}\right) f_{\mid i} P_{\mu \mu \lambda}^{n} \\
& +\sum_{i, j, k}\left(h_{k j}^{n+1} \Psi_{k j i}-h_{k i}^{n+1} \Psi_{k j j}\right) f_{\mid i} .
\end{aligned}
$$

By the second formula of Proposition 2.3, we can obtain that

$$
\sum_{j}\left(n H \delta_{i j}-h_{i j}^{n+1}\right)_{\mid j}-\sum_{j, k}\left(h_{k j}^{n+1} \Psi_{k j i}-h_{k i}^{n+1} \Psi_{k j j}\right)=0
$$

which together with Proposition 3.5 implies that the operator $\square$ is self-adjoint. When $\mathrm{H}_{2}$ is constant, we have the following computation by (4.3)

$$
\begin{align*}
\square(n H)= & \sum_{i} n H(n H)_{|i| i}-\sum_{i, j, k} h_{i j}^{n+1} h_{k k|i| j}^{n+1} \\
& +\sum_{i, \lambda, \mu}\left(n H \delta_{i \lambda}-h_{i \lambda}^{n+1}\right)(n H)_{\mid i} P_{\mu \mu \lambda}^{n} \\
& +\sum_{i, j, k}\left(h_{k j}^{n+1} \Psi_{k j i}-h_{k i}^{n+1} \Psi_{k j j}\right)(n H)_{\mid i} \\
\geq & \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}-\sum_{i}\left(n H_{\mid i}\right)^{2} \\
& +\frac{(n-1)}{n}\left[S-n H_{2}\right]\left\{2(n-1) H_{2}-\frac{n-2}{n} S\right.  \tag{4.4}\\
& +\sum_{i, \lambda, \mu}\left(n H \delta_{i \lambda}-h_{i \lambda}^{n+1}\right)(n H)_{\mid i} P_{\mu \mu \lambda}^{n} \\
& +\sum_{i, j, k}\left(h_{k j}^{n+1} \Psi_{k j i}-h_{k i}^{n+1} \Psi_{k j j}\right)(n H)_{\mid i} .
\end{align*}
$$

Using the fact that $P_{i j \lambda}^{n}=-A_{i j \lambda \mid n}$, by (3.12) and (3.8) we have

$$
\begin{align*}
\square(n H) \geq & \sum_{i, j, k}\left(h_{i j \mid k}^{n+1}\right)^{2}-\sum_{i}\left(n H_{\mid i}\right)^{2} \\
& +\frac{(n-1)}{n} \tag{4.5}
\end{align*} \quad\left[S-n H_{2}\right]\left\{2(n-1) H_{2}-\frac{n-2}{n} S\right\}
$$

On the other hand, let $x=\bar{x}^{a} \frac{\partial}{\partial \bar{x}^{a}}$ be the position vector field of the Minkowski space $V^{n+1}$ with respect to the origin. By a direct simple computation, we get
$\bar{\nabla}_{Z} x=Z, \quad \forall Z=z^{a} \frac{\partial}{\partial \bar{x}^{a}}$ on $V^{n+1}$. This together with the second formula of Proposition 2.1 implies that $\nabla_{e_{i}} x^{2}=2\left\langle e_{i}, x\right\rangle$ and $\nabla_{e_{i}}\left\langle e_{i}, x\right\rangle=\theta_{i}^{j}\left(e_{i}\right)\left\langle e_{j}, x\right\rangle+$ $h_{i i}^{n+1}\left\langle e_{n+1}, x\right\rangle+1$. Then, when $M$ is compact, there exists a point $P \in M$ such that $h_{i i}^{n+1}(P)>0, \forall i$, hence we have that $n(n-1) H_{2}(P)=\sum_{1 \leq i<j \leq n} \lambda_{i}(P) \lambda_{j}(P)>$ 0 , thus the constant $H_{2}>0$, which yields $(n H)^{2}>S$. On the other hand, we have

$$
\begin{equation*}
\sum_{i}(n H)^{2}\left(n H_{\mid i}\right)^{2}=\sum_{i}\left[\sum_{j, k} h_{j k}^{n+1} h_{j k \mid i}^{n+1}\right]^{2} \leq S \sum_{i, j, k}\left[h_{j k \mid i}^{n+1}\right]^{2} \tag{4.6}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\left(n H_{\mid i}\right)^{2} \leq\left(h_{j k \mid i}^{n+1}\right)^{2} \tag{4.7}
\end{equation*}
$$

It can be seen that our assumption $S \leq \frac{n(n-1)}{n-2} H_{2}$ is equivalent to

$$
\begin{equation*}
2(n-1) H_{2}-\frac{n-2}{n} S-\frac{n-2}{n} \sqrt{\left(n(n-1) H_{2}+S\right)\left(S-n H_{2}\right)} \geq 0 \tag{4.8}
\end{equation*}
$$

therefore the right-hand side of (4.5) is non-negative by (4.7) and (4.8). Because of the compactness of $M$, we get that $h_{i j}^{n+1}$ is constant and $h_{i j}^{n+1}=0, \forall i \neq j$ on $M$. Exterior differentiation of $h_{n a}^{n+1}=0$ yields $h_{a a}^{n+1}=h_{n n}^{n+1}, \quad \forall a=1, \cdots, n-1$, i.e., $h_{i i}^{n+1}=H, \forall i$. Since $h_{i j \mid k}^{n+1}=0$, it can be seen from the second formula of Proposition 2.3 that

$$
\begin{equation*}
h_{l k}^{n+1} \Psi_{l i j}-h_{l j}^{n+1} \Psi_{l i k}=0 \tag{4.9}
\end{equation*}
$$

Set $j=n, k=\lambda$ in (4.9); by $h_{i j}^{n+1}=0, \forall i \neq j$ we obtain that

$$
\begin{equation*}
h_{\lambda \lambda}^{n+1} \Psi_{\lambda i n}-h_{n n}^{n+1} \Psi_{n i \lambda}=0, \tag{4.10}
\end{equation*}
$$

which together with (2.4) yields

$$
\begin{equation*}
h_{n n}^{n+1} \bar{A}_{i j n+1}=0, \forall i, j \tag{4.11}
\end{equation*}
$$

So we get that $\Psi_{i j k}=0$ on $M$. Using the first formula of Proposition 2.2, we have that $P_{i j \lambda}^{n}=0$ on $M$, thus $M$ is a Landsberg space. It is easy to see from the second formula of Proposition 2.2 and $h_{i j}^{n+1}=H \delta_{i j}$ that

$$
\begin{equation*}
R_{i k l}^{j}=H\left(\delta_{i l} \delta_{j k}-\delta_{i k} \delta_{j l}\right) \tag{4.12}
\end{equation*}
$$

It can be seen from $h_{i i}^{n+1}=H$ and the constant $H_{2}=\frac{1}{n(n-1)}\left[(n H)^{2}-S\right]>0$ that $H \neq 0$ is constant, hence we get that $M$ is a Landsberg space with the nonzero constant flag curvature $H$, which together with Lemma 3.4 finishes the proof of Theorem 4.2.

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